

Sewing string tree vertices with ghosts

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Abstract

It is shown how to sew string vertices with ghosts at tree level in order to produce new tree vertices using the Group Theoretic approach to String Theory. It is then verified the BRST invariance of the sewn vertex and shown that it has the correct ghost number.

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1 Introduction

In the early days of String Theory, one way to obtain amplitudes for the scattering of an arbitrary number of strings was by using the factorization property, what means that the scattering amplitude of N strings may be interpreted as the scattering amplitudes of a smaller number of strings sewn together. This made it possible to build the N string scattering amplitude by knowing the expression for three string scattering amplitudes. Even though they were very ingenious and successful, those calculations didn't take into account the ghost structure of the vertices, and that is what is done here.

In [1], it was shown how to sew tree vertices without ghosts using the Group Theoretic approach to String Theory [2] in order to obtain a new, composite vertex. Following the same procedure, we shall calculate the scattering amplitude of N strings taking account the ghost structure.

We shall start with a short review of how to sew tree vertices without ghosts. What we must do is sew two legs of two vertices, one leg from each vertex. What we have in the beginning are two vertices V_1 and V_2 with N_1 and N_2 legs, respectively (figure 1).

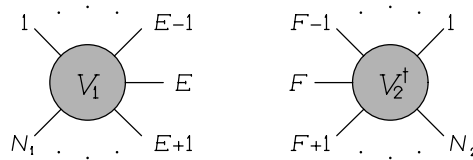


Figure 1: Individual vertices.

We now sew leg E from V_1 with the adjoint of leg F from V_2 . What we have now is the substitution of the two sewn legs by a propagator (figure 2). When this propagator is written in parametric form, it is an integration of one of the variables (in order to cancel one spurious degree of freedom) and a conformal factor \mathcal{P} which contains terms of L_n 's acting on leg E only.

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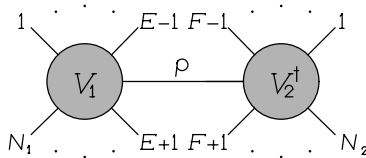


Figure 2: Sewn vertices.

So the resulting vertex V_c (called the *composite vertex*) has the generic form

$$V_c = V_1 P V_2^\dagger, \quad (1)$$

where the hermitian conjugate of V_2 is for the sewn leg F only and

$$P = \int dx \mathcal{P} \quad (2)$$

where x is a suitable variable. In what follows, we shall often write \mathcal{P} instead of P , calling attention to the integration when necessary.

When the two vertices are sewn together, we identify legs E and F . We also identify the Koba-Nielsen variable z_E with one of the Koba-Nielsen variables of vertex V_2^\dagger , and the Koba-Nielsen variable z_F we identify with one of the Koba-Nielsen variables of vertex V_1 . In [1], this identification is made in the following way: z_E may be identified with z_{F-1} or z_{F+1} , and z_F may be identified with z_{E-1} or z_{E+1} . So there are four possible combinations: a) $z_E = z_{F-1}$, $z_F = z_{E-1}$; b) $z_E = z_{F+1}$, $z_F = z_{E+1}$; c) $z_E = z_{F-1}$, $z_F = z_{E+1}$; d) $z_E = z_{F+1}$, $z_F = z_{E-1}$.

2 Oscillator case

The two original vertices satisfy some overlap identities and so shall do the composite vertex. One particular overlap identity is given by considering the operator $Q^{\mu i}$ with conformal weight $d = 0$ defined by [2]

$$Q^{\mu i}(\xi_i) = - \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} \frac{1}{\sqrt{n}} \alpha_n^{\mu i}(\xi_i)^{-n} + \alpha_0^{\mu i} \ln \xi_i + \frac{\bar{\partial}}{\partial \alpha_{0\mu}^i}. \quad (3)$$

where $\alpha_n^{\mu i}$ are bosonic oscillators with commutation relations

$$[\alpha_n^{\mu i}, a_m^{\nu j}] = -\eta^{\mu\nu} \delta^{ij} \delta_{n,-m}, \quad n, m \neq 0, \quad (4)$$

$$[\alpha_0^{\mu i}, a_n^{\nu j}] = 0, \quad \forall n. \quad (5)$$

The overlap identity is given by

$$V [Q^{\mu i}(\xi_i) - Q^{\mu j}(\xi_j)] = 0. \quad (6)$$

Because we are considering the adjoint of leg F in vertex V_2 , we must see what the adjoint of these overlap identities are. First, by the definition of $Q^{\mu i}$, we have that

$$Q^{\mu i \dagger}(\xi_i) = \Gamma Q^{\mu i}(\xi_i) \Gamma = Q^{\mu i}(\xi_i). \quad (7)$$

So, the adjoint of the overlap equations is given by

$$[Q^{\mu i}(\Gamma \xi_i) - Q^{\mu j}(\xi_j)] V^\dagger = 0. \quad (8)$$

We now take the overlap identity considering the effects of the operator $Q^{\mu i}(\xi_i)$ on the vertex V_1 on a generic leg i and on leg E^1 (figure 3):

¹As we shall be seeing soon, this form of the overlap will not lead to the correct composite vertex in the case where the cycling transformations of the legs that are not sewn involve the sewn legs E or F .

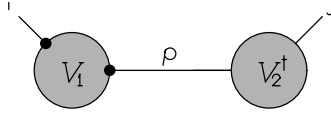


Figure 3: Overlap identity for V_1 .

$$V_1 \left[Q^{\mu i}(\xi_i) - Q^{\mu E}(\xi_E) \right] = 0 . \quad (9)$$

We may then insert the unit operator $1 = \mathcal{P}\mathcal{P}^{-1}$ and multiply by \mathcal{P} without altering the result:

$$V_1 \mathcal{P} \mathcal{P}^{-1} \left[Q^{\mu i}(\xi_i) - Q^{\mu E}(\xi_E) \right] \mathcal{P} . \quad (10)$$

Since the conformal operator \mathcal{P} acts only on leg E , it will have no effect on $Q^{\mu i}(\xi_i)$. In order to compute the effect of \mathcal{P} on $Q^{\mu E}(\xi_E)$, we must know that, for a conformal transformation V acting on a conformal operator $R(z)$ of weight d ,

$$V R(z) V^{-1} = \left(\frac{dVz}{dz} \right)^d R(z) . \quad (11)$$

Since $Q^{\mu E}(\xi_E)$ has conformal weight $d = 0$, we have

$$\mathcal{P}^{-1} Q^{\mu E}(\xi_E) \mathcal{P} = Q^{\mu E}(\mathcal{P}^{-1} \xi_E) , \quad (12)$$

and (figure 4)

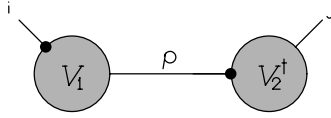


Figure 4: Overlap identity for $V_1 \mathcal{P}$.

$$V_1 \mathcal{P} \left[Q^{\mu i} \xi_i - Q^{\mu E}(\mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (13)$$

The second term in the expression above is facing now leg F of vertex V_2^\dagger , or best, its Hermitian conjugate. Considering that the Hermitian conjugate of $Q^{\mu i}$ (given by (7)), we then have the following overlap identity between legs i and F :

$$V_1 \mathcal{P} \left[Q^{\mu i}(\xi_i) - Q^{\mu F}(\Gamma \mathcal{P}^{-1} \xi_F) \right] = 0 . \quad (14)$$

We can then make a cycling transformation in order to obtain the correct factor for an arbitrary leg j ($j \neq F$) of vertex V_2^\dagger . The only term that will be affected is the term depending on leg F :

$$V_j^{-1} V_F Q^{\mu F}(\Gamma \mathcal{P}^{-1} \xi_E) V_F^{-1} V_j = Q^{\mu j} \left(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E \right) . \quad (15)$$

Doing this, the overlap identity for the composite vertex V_c (figure 5) can be written as

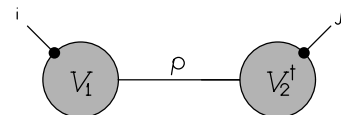


Figure 5: Overlap identity for V_c .

$$V_c \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 , \quad (16)$$

which is the overlap equation between two arbitrary legs i and j of the composite vertex V_c .

But the overlap identity for the composite vertex, since none of legs i or j involves the propagator, must be given by

$$V_c \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(\xi_j) \right] = 0 \quad (17)$$

and so in order for the equation we have obtained for the overlap of the composite vertex V_c to be true we must have

$$V_F \Gamma \mathcal{P}^{-1} V_E^{-1} = 1 \Rightarrow \mathcal{P}^{-1} = \Gamma^{-1} V_F^{-1} V_E \quad (18)$$

which implies that the propagator is given by

$$\mathcal{P} = V_E^{-1} V_F \Gamma . \quad (19)$$

In order to give an explicit expression for the propagator, we will now choose ξ_i to be of the form

$$\xi_i = V_i^{-1} z = z - z_i . \quad (20)$$

This choice is called the ‘‘simple cycling’’ [1] and it is the one that simplifies our calculations the most. In this choice, the propagator is given by

$$\mathcal{P} z = \frac{1}{z} + z_F - z_E , \quad (21)$$

or in terms of the L_n^E operators²,

$$\mathcal{P} = e^{(z_F - z_E - 1)L_{-1}^E} (-1)^{L_0^E} e^{L_1^E} e^{-L_{-1}^E} . \quad (22)$$

This form works for all choices for the composite vertex discussed before³. The true propagator is given by expression (22) integrated over a suitable variable. Choosing this variable to be $s = z_F - z_E$, we then have

$$P = \int_{-\infty}^0 ds \mathcal{P} = \frac{1}{L_{-1}^E} (-1)^{L_0^E} e^{L_1^E} e^{-L_{-1}^E} , \quad s = z_F - z_E . \quad (23)$$

Before going any further, we must discuss another aspect of the theory that depends on the particular way in which the legs are identified during the sewing procedure. Let us consider the more general case of an arbitrary cycling V_i . This kind of cycling may depend on other coordinates that are not z_i . As an example, let us suppose that we are identifying coordinate z_E of leg E with coordinate z_{F-1} of vertex V_2^\dagger and coordinate z_F of leg F with coordinate z_{F-1} of vertex V_1 . The overlap identity between legs i and $E-1$ on vertex V_1 is

$$V_1 \left[Q^{\mu i} \left(V_{0i}^{-1} z \right) - Q^{\mu(E-1)} \left(V_{0(E-1)}^{-1} z \right) \right] \quad (24)$$

where we are calling V_{0j}^{-1} the cycling transformation on leg j ($j = 1, \dots, N$). These cycling transformations may depend on the other legs. As an example when this happens, we take another choice of the cycling transformations ξ_i that is not as trivial as (20) but gives a simpler formula for the propagator. This choice is given by [1] [3]

$$\xi_i = V_{0i}^{-1} z = \frac{(z_{i+1} - z_{i-1})}{(z_{i+1} - z_i)} \cdot \frac{(z - z_i)}{(z - z_{i-1})} \quad (25)$$

²Other forms for this propagator are given by [1]:

$$\begin{aligned} \mathcal{P} &= e^{-L_1^E/s} (-1)^{L_0^E} s^{2L_0^E} e^{L_{-1}^E/s} \\ \mathcal{P} &= (-1)^{L_0^E - L_1^E/s} s^{2(L_0^E - L_1^E/s)} e^{-(L_1^E - L_{-1}^E)/s} . \end{aligned}$$

³This affirmation is usually not valid for other choices of ξ_i .

which is the transformation that takes z_{i-1} , z_i and z_{i+1} to ∞ , 0 and 1, respectively. Its inverse is given by

$$V_{0i}z = \frac{z_{i-1}(z_i - z_{i+1})z + z_i(z_{i+1} - z_{i-1})}{(z_i - z_{i+1})z + (z_{i+1} - z_{i-1})} . \quad (26)$$

In the case of the cycling given by (25), the cycling for leg $E - 1$ will depend on leg E , which is not present in the composite vertex. On this vertex, the overlap between leg $E - 1$ and an arbitrary leg j reads

$$V_c \left[Q^{\mu(E-1)} \left(V_{0(E-1)}^{-1} z \right) - Q^{\mu j} \left(V_{0j}^{-1} z \right) \right] . \quad (27)$$

This overlap equation involves terms that depend on leg E , which is non-existent in the composite vertex V_c . The correct overlaps should be given by

$$V_c \left[Q^{\mu(E-1)} \left(V_{E-1}^{-1} z \right) - Q^{\mu j} \left(V_j^{-1} z \right) \right] , \quad (28)$$

where the cycling transformations V_{E-1}^{-1} and V_j^{-1} do not depend on legs E or F . So, in order to restore the correct cycling transformation for the composite vertex, a conformal transformation must be made on the cycling transformations on vertex V_1 . These are given by

$$V_{0j}^{-1} \longrightarrow C_1 V_{0j}^{-1} , \quad j \neq E , \quad (29)$$

where

$$C_1 = \prod_{i=1}^{N_1} V_i^{-1} V_{0i} . \quad (30)$$

In this definition, we consider implicit that the transformation $V_E^{-1} V_{0E} = 1$ since the conformal transformations on leg E will not be part of the composite vertex and so need not be modified.

Considering the general case, we have that the overlap equation for V_1 (figure 6) that will lead to the correct composite vertex is now obtained from the original overlap

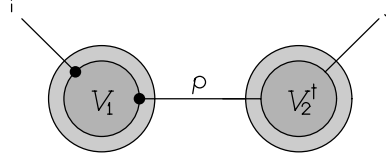


Figure 6: Overlap identity for V_1 .

$$V_1 \left[Q^{\mu i}(\xi_{0i}) - Q^{\mu E}(\xi_{0E}) \right] = 0 , \quad (31)$$

where $\xi_{0i} = V_{0i}^{-1} z$ and $\xi_{0E} = V_{0E}^{-1} z$. By inserting conformal transformation (30), we obtain (figure 7)

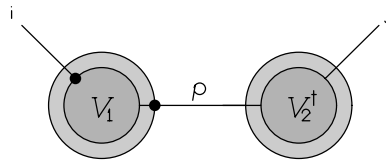


Figure 7: Overlap identity for $V_1 C_1^{-1}$.

$$V_1 C_1^{-1} \left[Q^{\mu i}(\xi_i) - Q^{\mu E}(\xi_E) \right] = 0 . \quad (32)$$

Inserting now the propagator, we obtain (figure 8)

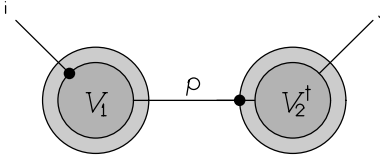


Figure 8: Overlap identity for $V_1 C_1^{-1} \mathcal{P}$.

$$V_1 C_1^{-1} \mathcal{P} \left[Q^{\mu i}(\xi_i) - Q^{\mu E}(\mathcal{P}^{-1} \xi_E) \right] = 0 \quad (33)$$

and we expect the composite vertex to have a different form (given shortly) than in (1) in order to amount for the contributions of the conformal transformations. The second term of the overlap is now facing leg F of vertex V_2^\dagger so that we have the following overlap between legs i and F :

$$V_1 C_1^{-1} \mathcal{P} \left[Q^{\mu i}(\xi_i) - Q^{\mu F}(\Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (34)$$

We are now facing the conformal transformation C_F^4 , defined by

$$C_F = V_F^{-1} V_{0F} \quad (35)$$

which is necessary in order to change $\xi_{0F} \rightarrow \xi_F$. Introducing this transformation we obtain (figure 9)

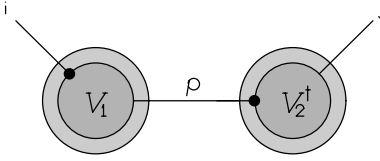


Figure 9: Overlap identity for $V_1 C_1^{-1} \mathcal{P} C_F$.

$$V_1 C_1^{-1} \mathcal{P} C_F \left[Q^{\mu i}(\xi_i) - Q^{\mu F}(V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (36)$$

Making a cycling transformation from leg F to leg j , we then obtain (figure 10)

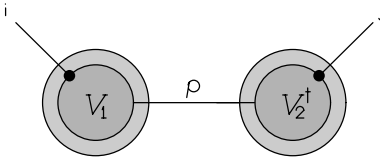


Figure 10: Overlap identity for $V_1 C_1^{-1} \mathcal{P} C_F V_2^\dagger$.

$$V_1 C_1^{-1} \mathcal{P} C_F V_2^\dagger \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(V_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (37)$$

Once again, a conformal transformation must be introduced because of the cycling transformations V_{0j}^{-1} . This is defined by

$$C_2 = \prod_{\substack{i=1 \\ i \neq F}}^{N_2} V_i^{-1} V_{0i} . \quad (38)$$

so that we now have (figure 11)

⁴In this case, like we have seen for leg V_E , V_F may depend on the variable z_E .

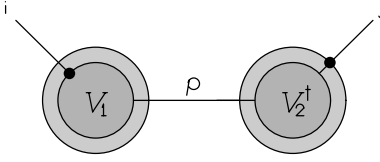


Figure 11: Overlap identity for $V_1 C_1^{-1} \mathcal{P} C_F V_2^\dagger C_2^{-1}$.

$$V_1 C_1^{-1} \mathcal{P} C_F V_2^\dagger C_2^{-1} \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (39)$$

The composite vertex must be defined in terms of the new cycling transformations and so it must now include the conformal transformations that perform this change. So, it will now be defined by

$$V_c = V_1 C_1^{-1} \mathcal{P} C_F V_2^\dagger C_2^{-1} . \quad (40)$$

Considering this, the overlap identity for the composite vertex V_c can be written as

$$V_c \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0 . \quad (41)$$

Since the correct overlap identity for the composite vertex is given by

$$V_c \left[Q^{\mu i}(\xi_i) - Q^{\mu j}(\xi_j) \right] = 0 \quad (42)$$

we must have

$$V_F \Gamma \mathcal{P}^{-1} V_E^{-1} = 1 \quad (43)$$

which implies once again that

$$\mathcal{P} = V_E^{-1} V_F \Gamma . \quad (44)$$

For the cycling transformation (25), it is only necessary to do conformal transformations on legs $E-1$, E , $E+1$, $F-1$, F and $F+1$, depending on the particular way the variables associated with these legs are identified with the variables associated to legs E and F . In this particular example (which is case a seen before), the conformal transformations are given by

$$V_{0(E-1)}^{-1} \longrightarrow C_1 V_{0(E-1)}^{-1} \quad , \quad C_1 = r^{L_0^{E-1}} , \quad (45)$$

$$V_{0(F-1)}^{-1} \longrightarrow C_2 V_{0(F-1)}^{-1} \quad , \quad C_2 = t^{L_0^{F-1}} \quad (46)$$

where

$$r = \frac{(z_{F-1} - z_{E-2})(z_{E-1} - z_E)}{(z_{F-1} - z_{E-1})(z_{E-2} - z_E)} \quad , \quad t = \frac{(z_{E-1} - z_{F-2})(z_{F-1} - z_F)}{(z_{E-1} - z_{F-1})(z_{F-2} - z_F)} . \quad (47)$$

For the cycling (25), the propagators obtained for the four possible combinations discussed before are given by

$$\text{a) } \mathcal{P}_a = a^{L_0^E} \quad , \quad \text{b) } \mathcal{P}_b = e^{-L_1^E} a^{L_0^E} e^{-L_{-1}^E} \quad , \quad \text{c) } \mathcal{P}_c = e^{-L_1^E} b^{L_0^E} \quad , \quad \text{d) } \mathcal{P}_d = b^{L_0^E} e^{-L_{-1}^E} , \quad (48)$$

where⁵

$$a = \frac{(z_{E+1} - z_{E-1})(z_{F+1} - z_{F-1})}{(z_{E+1} - z_{F-1})(z_{F+1} - z_{E-1})} \quad , \quad b = \frac{(z_{E+1} - z_{E-1})(z_{F+1} - z_{F-1})}{(z_{E+1} - z_{F+1})(z_{F-1} - z_{E-1})} . \quad (49)$$

The true propagators are obtained when we integrate the expressions above multiplied by a suitable constant. The results are:

$$\text{a) } P_a = \int_0^1 da a^{L_0^E - 1} = \frac{1}{L_0^E} \quad , \quad \text{b) } P_b = \int_0^1 da e^{-L_1^E} a^{L_0^E - 1} e^{-L_{-1}^E} = e^{-L_1^E} \frac{1}{L_0^E} e^{-L_{-1}^E} , \quad (50)$$

$$\text{c) } P_c = \int_0^1 db e^{-L_1^E} b^{L_0^E - 1} = e^{-L_1^E} \frac{1}{L_0^E} \quad , \quad \text{d) } P_d = \int_0^1 db b^{L_0^E - 1} e^{-L_{-1}^E} = \frac{1}{L_0^E} e^{-L_{-1}^E} . \quad (51)$$

⁵The coefficient a can be connected with the coefficient c in reference [1] by $a = \frac{c}{c-1}$.

It is now necessary to verify the effect of the gauge transformations C_1 and C_2 on the composite vertex as given by formula (40). We shall do it by verifying the effect of C_1 on vertex V_1 . In order to do this we need the explicit expression for the bosonic oscillator vertex V_1 , given by [1] [4]

$$V_1 = \left(\prod_{i=1}^N \langle 0|_i \right) \exp \left[-\frac{1}{2} \sum_{\substack{i,j=1 \\ i \neq j}}^{N_1} \sum_{n,m=0}^{\infty} \alpha_n^{\mu i} D_{nm} \left(\Gamma V_{0i}^{-1} V_{0j} \right) \alpha_{m\mu}^j \right] \quad (52)$$

where V_{0i}^{-1} and V_{0j} are cycling transformations involving leg E and the oscillators $\alpha_n^{\mu i}$ have commutation relations given by (4, 5). Matrices $D_{nm}(\gamma)$ are defined in the following way [3]:

$$D_{n0}(\gamma) = \frac{1}{\sqrt{n}} [\gamma(0)]^n, \quad (53)$$

$$D_{nm}(\gamma) = \sqrt{\frac{m}{n}} \frac{1}{m!} \frac{\partial^m}{\partial z^m} [\gamma(z)]^n \Big|_{z=0}, \quad (54)$$

$$D_{00}(\gamma) = \frac{1}{2} \ln \left[\frac{d}{dz} \gamma(z) \right] \Big|_{z=0} \quad (55)$$

and have the following multiplication property:

$$D_{nm}(\gamma_1 \gamma_2) = \sum_{p=1}^{\infty} D_{np}(\gamma_1) D_{pm}(\gamma_2) + D_{n0}(\gamma_1) \delta_{0m} + \delta_{0n} D_{0m}(\gamma_2). \quad (56)$$

In order to calculate the effects of the conformal transformation C_1 on these oscillators we must make use of the following conformal operator [2]

$$P^{\mu i}(\xi_{0i}) = \sum_{\substack{n=-\infty \\ n \neq 0}}^{\infty} \sqrt{|n|} \alpha_n^{\mu i} (\xi_{0i})^{-n-1} + \alpha_0^{\mu i} (\xi_{0i})^{-1} \quad (57)$$

which has conformal weight one, what means it transforms like

$$P^{\mu i}(\xi_{0i}) = \frac{d\xi_{0j}}{d\xi_{0i}} P^{\mu j}(\xi_{0j}). \quad (58)$$

An oscillator $\alpha_n^{\mu i}$ ($n \geq 1$) can be expressed in terms of this conformal operator in the following way:

$$\alpha_n^{\mu i} = \frac{1}{\sqrt{n}} \oint_{\xi_{0i}=0} d\xi_{0i} (\xi_{0i})^n P^{\mu i}(\xi_{0i}). \quad (59)$$

Acting on it with the conformal transformation C_1 , we have

$$C_1 \alpha_n^{\mu i} C_1^{-1} = \frac{1}{\sqrt{n}} \oint_{\xi_{0i}=0} d\xi_{0i} (\xi_{0i})^n \left(\frac{d}{d\xi_{0i}} V_i^{-1} V_{0i} \xi_{0i} \right) P^{\mu i}(V_i^{-1} V_{0i} \xi_{0i}). \quad (60)$$

Making now a change of variables $\xi_{0i} \rightarrow \xi_i$, we have

$$C_1 \alpha_n^{\mu i} C_1^{-1} = \frac{1}{\sqrt{n}} \oint_{\xi_i=0} d\xi_i (V_{0i}^{-1} V_i \xi_i)^n P^{\mu i}(\xi_i). \quad (61)$$

Expanding $(V_{0i}^{-1} V_i \xi_i)^n$ in terms of $D_{nm}(\gamma)$ matrices and $P^{\mu i}(\xi_i)$ in terms of the oscillators, we obtain

$$\begin{aligned} C_1 \alpha_n^{\mu i} C_1^{-1} &= \\ & \sum_{\substack{p=-\infty \\ p \neq 0}}^{\infty} \sqrt{\frac{p}{n}} \oint_{\xi_i=0} d\xi_i \left[\sqrt{n} D_{n0} (V_{0i}^{-1} V_i) (\xi_i)^{-p-1} + \sum_{m=1}^{\infty} \sqrt{\frac{n}{m}} D_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m-p-1} \right] \alpha_p^{\mu i} \\ & + \frac{1}{\sqrt{n}} \oint_{\xi_i=0} d\xi_i \left[\sqrt{n} D_{n0} (V_{0i}^{-1} V_i) (\xi_i)^{-1} + \sum_{m=1}^{\infty} \sqrt{\frac{n}{m}} D_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m-1} \right] \alpha_0^{\mu i}. \end{aligned} \quad (62)$$

Performing the integrations we then obtain

$$C_1 \alpha_n^{\mu i} C_1^{-1} = \sum_{m=0}^{\infty} D_{nm} \left(V_{0i}^{-1} V_i \right) \alpha_m^{\mu i} . \quad (63)$$

Using the same process for $\alpha_0^{\mu i}$, we obtain

$$C_1 \alpha_0^{\mu i} C_1^{-1} = \alpha_0^{\mu i} . \quad (64)$$

Using these transformation properties, the multiplication rules of matrices $D_{nm}(\gamma)$ (equation (56)) and the property [3]

$$D_{nm}(\gamma) = D_{mn}(\Gamma \gamma^{-1} \Gamma) , \quad (65)$$

we can show that the effect of the conformal transformation C_1 on vertex V_1 is given by

$$V_1 C_1^{-1} = \left(\prod_{i=1}^N \langle 0 | \right) \exp \left[-\frac{1}{2} \sum_{\substack{i,j=1 \\ i \neq j}}^{N_1} \sum_{n,m=0}^{\infty} \alpha_n^{\mu i} D_{nm} \left(\Gamma V_i^{-1} V_j \right) \alpha_{m\mu}^j \right] , \quad (66)$$

i.e. the effect of C_1 on vertex V_1 is to change $V_{0i}^{-1} \rightarrow V_i^{-1}$ and $V_{0j} \rightarrow V_j$ thus eliminating the dependence of the cycling transformations of every leg except leg E on the latter. The same can be done to obtain the effect of conformal transformation C_2 on vertex V_2^\dagger , with the same results. So, the effect of these transformations is to eliminate from the cycling transformations of the composite vertex all dependence on the sewn legs E and F .

3 Introduction of ghosts

We now introduce ghosts in the vertex so that what we must sew now are two vertices with some ghost variables attached to them, i.e. we will be considering the physical vertices [4] which have the correct ghost number. In this case, in addition to satisfying the overlap identities with the conformal operator $Q^{\mu i}$, the two physical vertices and the composite vertex must also satisfy some overlap identities with the conformal operators b^i and c^i , given by [4]

$$b^i(\xi_i) = \sum_{n=-\infty}^{\infty} b_{-n}^i(\xi_i)^{n-2} , \quad (67)$$

$$c^i(\xi_i) = \sum_{n=-\infty}^{\infty} c_{-n}^i(\xi_i)^{n+1} \quad (68)$$

where c_n^i and b_n^i are ghost anticommuting oscillators with anticommutation relations

$$\{c_n^i, b_m^j\} = \delta_{n,-m} . \quad (69)$$

These operators have, respectively, conformal weights 2 and -1 , what means that they transform like

$$b^i(\xi_i) = \left(\frac{d\xi_j}{d\xi_i} \right)^2 b^j(\xi_j) , \quad (70)$$

$$c^i(\xi_i) = \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} c^j(\xi_j) . \quad (71)$$

The overlap identities for a vertex V with these operators are given by

$$V \left[b^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^2 b^j(\xi_j) \right] = 0 , \quad (72)$$

$$V \left[c^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} c^j(\xi_j) \right] = 0 . \quad (73)$$

We shall be working here with overlap identities for the physical vertex U [4], which has the correct ghost number, instead of the overlap identities for vertex V . The physical vertex is given by [4]

$$U = V \prod_{\substack{i=1 \\ i \neq a, b, c}}^N \sum_{j=1}^N \sum_{n=-1}^{\infty} e_n^{ij} b_n^j \quad (74)$$

where a, b, c are any three legs of the vertex and the matrix e_n^{ij} is given by

$$\sum_{n=-1}^{\infty} e_n^{ij} \mathcal{L}_n^j = V_j^{-1} \partial_{z_i} V_j \quad (75)$$

where the cycling transformations are now defined on the complete generators \mathcal{L}_n^i of the conformal algebra of the bosonic oscillators and of the ghost oscillators. These vectors e_n^{ij} have the following property:

$$\frac{\partial V}{\partial z_i} = V \sum_{j=1}^N \sum_{n=-1}^{\infty} e_n^{ij} \mathcal{L}_n^j . \quad (76)$$

In order to derive the overlap identity for the physical vertex U , we must multiply the overlap identity for V by the same factor as in equation (74),

$$V \left[b^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^2 b^j(\xi_j) \right] \left(\prod_{\substack{k=1 \\ k \neq a, b, c}}^N \sum_{l=1}^N \sum_{n=-1}^{\infty} e_n^{kl} b_n^l \right) = 0 , \quad (77)$$

$$V \left[c^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} c^j(\xi_j) \right] \left(\prod_{\substack{k=1 \\ k \neq a, b, c}}^N \sum_{l=1}^N \sum_{n=-1}^{\infty} e_n^{kl} b_n^l \right) = 0 \quad (78)$$

and pass it through the overlap identities, obtaining

$$U \left[b^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^2 b^j(\xi_j) \right] = 0 , \quad (79)$$

$$U \left[c^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} c^j(\xi_j) \right] + V \sum_{\substack{p=1 \\ p \neq a, b, c}}^N (-1)^p \prod_{\substack{k=1 \\ k \neq a, b, c \\ k \neq p}}^N \sum_{l=1}^N \sum_{q=-1}^{\infty} e_q^{kl} b_q^l \sum_{n=-1}^{\infty} \left[e_n^{pi}(\xi_i)^{n+1} - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} e_n^{pj}(\xi_j)^{n+1} \right] = 0 . \quad (80)$$

From (80) we can see that there will be an anomalous term in the c^i overlap of the physical vertex U unless both legs i and j are precisely those legs (a, b or c) that do not have any ghosts attached to them. These ghosts which are attached to all the other legs are responsible for the anomalous terms.

3.1 Analysis of the ghost number

Before going any further, it is necessary to make some considerations on the ghost number of the composite vertex. As we shall see shortly, in the case where we perform the sewing with ghosts included, using the physical vertices, the resulting composite physical vertex will not have the correct ghost number unless we insert some extra ghosts in vertex U_1 before the sewing takes place. Considering this, we shall define the composite vertex to be given by

$$U_c = U_1 G P U_2^\dagger \quad (81)$$

where G are some extra ghosts that will be introduced in order to make vertex U_c have the correct ghost number and P is the propagator (in its integrated form).

We must now analyze the ghost number of the composite vertex and of its parts in order to calculate the ghost number that the extra ghosts G must have. In order to do this, we shall use the ghost number operator N_{gh} . For a vertex with N legs, the ghost number operator is defined by

$$N^{\text{gh}} = \sum_{i=1}^N \left(\sum_{n=-1}^{\infty} c_{-n}^i b_n^i - \sum_{n=2}^{\infty} b_{-n}^i c_n^i \right). \quad (82)$$

The reason why the ghost number operator is a sum from $i = 1$ to $i = N$ is because there are N vacua that will annihilate the operators corresponding to each one of them. When acting on the physical vertex, this operator gives a ghost number N , what is the correct ghost number for a tree vertex with N legs.

In the case of the composite vertex, it has $N_1 + N_2 - 2$ legs (because it does not have legs E and F , which have been sewn together) and so it must have ghost number $(N_1 + N_2 - 2)$. For this vertex, the ghost number operator N_c^{gh} can be divided into two parts:

$$N_c^{\text{gh}} = N_1^{\text{gh}} + N_2^{\text{gh}} \quad (83)$$

where

$$N_1^{\text{gh}} = \sum_{\substack{i=1 \\ i \neq E}}^{N_1} \left(\sum_{n=-1}^{\infty} c_{-n}^i b_n^i - \sum_{n=2}^{\infty} b_{-n}^i c_n^i \right), \quad (84)$$

$$N_2^{\text{gh}} = \sum_{\substack{i=1 \\ i \neq F}}^{N_2} \left(\sum_{n=-1}^{\infty} c_{-n}^i b_n^i - \sum_{n=2}^{\infty} b_{-n}^i c_n^i \right). \quad (85)$$

This ghost number operator will have the following effect on the composite vertex:

$$U_c N_c^{\text{gh}} = U_c (N_1^{\text{gh}} + N_2^{\text{gh}}) = (N_1 + N_2 - 2) U_c. \quad (86)$$

Given formula (81) for the composite vertex, we then have

$$U_c N_c^{\text{gh}} = U_1 N_1^{\text{gh}} G P U_2^\dagger + U_1 [G, N_1^{\text{gh}}] P U_2^\dagger + U_1 G P U_2^\dagger N_2^{\text{gh}}. \quad (87)$$

In order to calculate this, we must pay some attention to terms one and three of the right hand side of the expression above. We know that

$$U_1 (N_1^{\text{gh}} + N_E^{\text{gh}}) = N_1 U_1, \quad (88)$$

$$U_2 (N_2^{\text{gh}} + N_F^{\text{gh}}) = N_2 U_2 \quad (89)$$

where

$$N_E^{\text{gh}} = \sum_{n=-1}^{\infty} c_{-n}^E b_n^E - \sum_{n=2}^{\infty} b_{-n}^E c_n^E, \quad (90)$$

$$N_F^{\text{gh}} = \sum_{n=-1}^{\infty} c_{-n}^F b_n^F - \sum_{n=2}^{\infty} b_{-n}^F c_n^F. \quad (91)$$

Taking the Hermitian conjugate of equation (89) only on leg F , we obtain

$$U_2^\dagger N_2^{\text{gh}} + N_F^{\text{gh}\dagger} U_2^\dagger = N_2 U_2^\dagger. \quad (92)$$

Since for $N_F^{\text{gh}\dagger}$ (and for any arbitrary ghost number operator)

$$N_F^{\text{gh}\dagger} = -N_F^{\text{gh}} + 3 \quad (93)$$

we then obtain, substituting (88), (92) and (93) into equation (87),

$$U_c N_c^{\text{gh}} = (N_1 + N_2 - 3)U_1 G P U_2^\dagger + U_1 \left[G, N_1^{\text{gh}} \right] P U_2^\dagger - U_1 N_E^{\text{gh}} G P U_2^\dagger + U_1 G P N_F^{\text{gh}} U_2^\dagger. \quad (94)$$

Passing N_E^{gh} through the extra ghosts G , we then obtain

$$\begin{aligned} U_c N_c^{\text{gh}} &= (N_1 + N_2 - 3)U_1 G P U_2^\dagger + U_1 \left[G, N_1^{\text{gh}} \right] P U_2^\dagger \\ &\quad - U_1 \left[N_E^{\text{gh}}, G \right] P U_2^\dagger - U_1 G P N_E^{\text{gh}} U_2^\dagger + U_1 G P N_F^{\text{gh}} U_2^\dagger. \end{aligned} \quad (95)$$

We must now remember that, in the composite vertex, we identify every operator on leg E with operators on leg F so that $N_E^{\text{gh}} = N_F^{\text{gh}}$. Doing this, the last two terms in (95) cancel and we obtain the following result:

$$U_c N_c^{\text{gh}} = (N_1 + N_2 - 3)U_1 G P U_2^\dagger + U_1 \left[G, N_1^{\text{gh}} + N_E^{\text{gh}} \right] P U_2^\dagger. \quad (96)$$

The fact that U_c has ghost number $N_1 + N_2 - 2$ then implies that

$$\left[G, N_1^{\text{gh}} + N_E^{\text{gh}} \right] = G, \quad (97)$$

i.e. the extra ghosts that must be introduced in vertex U_1 must have ghost number 1.⁶

There is an infinite number of combinations of ghosts that have ghost number 1. We could have any linear combination of ghosts of the type $b, bcb, bcbcb$, etc. but it will prove to be simpler to choose G to be a combination of b 's only so that we may represent it as

$$G = \sum_{i=1}^{N_1} \sum_{n=-\infty}^{\infty} \alpha_n^i b_n^i \quad (98)$$

where α_n^i are the coefficients of the linear combination. In order to determine the correct linear combination, we must use some other conditions, like BRST invariance of the scattering amplitude. This we shall see next.

3.2 BRST invariance

We must now impose that the scattering amplitude obtained from the composite vertex is BRST invariant and check whether this condition is strong enough to determine G . The scattering amplitude [1] is obtained by acting with the composite vertex

$$U_c = U_1 G P U_2^\dagger \quad (99)$$

on a certain number of physical states ($|\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle$) and then by integrating over all variables z_i ($i = 1, \dots, N_1 + N_2; i \neq E, F$):

$$W = \int \prod_{\substack{i=1 \\ i \neq E, F}}^{N_1+N_2} dz_i U_1 G P U_2^\dagger |\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle. \quad (100)$$

P is the propagator in its integrated form and G are the extra ghosts to be inserted in U_1 .

The action of the BRST charge Q on this scattering amplitude is given by

$$\begin{aligned} WQ &= \int \prod_{\substack{i=1 \\ i \neq E, F}}^{N_1+N_2} dz_i [U_1, Q] G P U_2^\dagger |\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle \\ &\quad + \int \prod_{\substack{i=1 \\ i \neq E, F}}^{N_1+N_2} dz_i U_1 [G, Q] P U_2^\dagger |\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle + \int \prod_{\substack{i=1 \\ i \neq E, F}}^{N_1+N_2} dz_i U_1 G [P, Q] U_2^\dagger |\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle \\ &\quad + \int \prod_{\substack{i=1 \\ i \neq E, F}}^{N_1+N_2} dz_i U_1 G P [U_2^\dagger, Q] |\chi_1\rangle|\chi_2\rangle \dots |\chi_N\rangle. \end{aligned} \quad (101)$$

⁶This contrasts with references [5] and [6] where it is claimed that the extra ghosts should have ghost number 1, 2 or 3, depending on the way one chooses legs E and F to have or not to have ghosts attached to them.

The first and third terms will result in total derivatives that give zero when one integrates over some variables [1] and so what remain are just the second and third terms.

The commutator P is given by a pure conformal transformation, and it is a function of the generators \mathcal{L}_n^E ($n = -1, \dots$) only. As the BRST charge commutes with all \mathcal{L}_n^E 's, i.e.

$$[\mathcal{L}_n^E, Q] = 0 \quad (102)$$

we have

$$[P, Q] = 0 . \quad (103)$$

Considering now that

$$[b_n^i, Q] = \mathcal{L}_n^i , \quad (104)$$

we then have, for G given by (98),

$$U_1[G, Q] = U_1 \sum_{i=1}^{N_1} \sum_{n=-1}^{\infty} \alpha_n^i \mathcal{L}_n^i . \quad (105)$$

In order for the scattering amplitude W to be BRST invariant, expression (105) must be zero or a total derivative (that can be integrated out to become a null surface term). At the same time, we want these extra ghosts to place (talking in terms of the simple cycling) a ghost on one of the legs in U_1 that do not have any ghosts attached to them. If we now remember property (76), we see that we can satisfy these constraints in a nice way by choosing G to be given by

$$G = (-1)^{N_1+a} \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{aj} b_n^j \quad (106)$$

where a ($a \neq E$) is one of the legs of vertex U_1 that does not have ghosts attached to it. Inserting these ghosts in vertex U_1 , we have

$$U_1 G = V_1 \prod_{\substack{i=1 \\ i \neq a, b, c}}^{N_1} \sum_{k=1}^{N_1} \sum_{m=-1}^{\infty} e_n^{ik} b_m^k \times (-1)^{N_1+a} \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{aj} b_n^j = V_1 \prod_{\substack{i=1 \\ i \neq b, c}}^{N_1} \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{ij} b_n^j . \quad (107)$$

Using formula (106) for the extra ghosts G , we then have

$$\begin{aligned} U_1[G, Q] &= U_1 (-1)^{N_1+a} \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{aj} [b_n^j, Q] = (-1)^{N_1+a} U_1 \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{aj} \mathcal{L}_n^j \\ &= (-1)^{N_1+a} \frac{\partial V_1}{\partial z_a} \prod_{\substack{i=1 \\ i \neq a, b, c}}^{N_1} \sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{ij} b_n^j \end{aligned} \quad (108)$$

what is a total derivative that will vanish when one does the integration over z_a .

In references [5] and [6], the extra ghosts have been placed in the propagator. Although this can be done, there is no way one can derive a formula for the ghosts in the propagator for a general cycling. In that case, the extra ghosts must be derived and BRST invariance has to be checked for each particular cycling. Also, the resulting composite vertex obtained in that case is not similar in its ghost structure to an ordinary tree vertex, although it has the correct ghost number.

3.3 Overlap identities

We must now use the overlap identities to determine the propagator that satisfies them. In order to do this we shall start with vertex U_1^0 , which is the vertex with cycling transformations V_i^{-1} which involve leg E . Considering equations (79) and (80), the overlap identities for vertex U_1^0 between an arbitrary leg i and leg E are given by (figure 12)

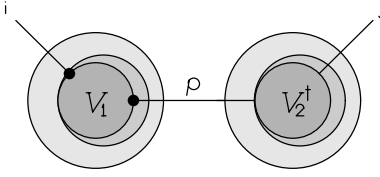


Figure 12: Overlap identity for U_1^0 .

$$U_1^0 \left[b^i(\xi_{0i}) - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^2 b^E(\xi_{0E}) \right] = 0 , \quad (109)$$

$$U_1^0 \left[c^i(\xi_{0i}) - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} c^E(\xi_{0E}) \right] + V_1^0 \sum_{\substack{p=1 \\ p \neq a, b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq a, b, c \\ k \neq p}}^{N_1} \sum_{l=1}^{N_1} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l \sum_{n=-1}^{\infty} \left[e_n^{pi}(\xi_{0i})^{n+1} - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} e_n^{pE}(\xi_{0E})^{n+1} \right] = 0 . \quad (110)$$

The extra ghosts must then be inserted in vertex U_1^0 so that the composite vertex will have the correct ghost number. Multiplying expressions (109) and (110) by the extra ghosts G (given by (106)) and passing them through the overlaps we obtain (figure 13)

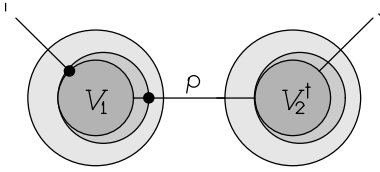


Figure 13: Overlap identity for $U_1^0 G$.

$$U_1^0 G \left[b^i(\xi_{0i}) - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^2 b^E(\xi_{0E}) \right] = 0 , \quad (111)$$

$$U_1^0 G \left[c^i(\xi_{0i}) - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} c^E(\xi_{0E}) \right] + U_1^0 (-1)^{N_1+a} \sum_{n=-1}^{\infty} \left[e_n^{ai}(\xi_{0i})^{n+1} - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} e_n^{aE}(\xi_{0E})^{n+1} \right] + V_1^0 \sum_{\substack{p=1 \\ p \neq a, b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq a, b, c \\ k \neq p}}^{N_1} \left(\sum_{l=1}^{N_1} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l \right) \left(\sum_{j=1}^{N_1} \sum_{n=-1}^{\infty} e_n^{aj} b_n^j \right) \times \sum_{n=-1}^{\infty} \left[e_n^{pi}(\xi_{0i})^{n+1} - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} e_n^{pE}(\xi_{0E})^{n+1} \right] = 0 . \quad (112)$$

The second and third terms of equation (112) can be combined so that it becomes

$$U_1^0 G \left[c^i(\xi_{0i}) - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} c^E(\xi_{0E}) \right] + V_1^0 \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \left(\sum_{l=1}^{N_1} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l \right) \sum_{n=-1}^{\infty} \left[e_n^{pi}(\xi_{0i})^{n+1} - \left(\frac{d\xi_{0E}}{d\xi_{0i}} \right)^{-1} e_n^{pE}(\xi_{0E})^{n+1} \right] = 0 . \quad (113)$$

At this point, we must introduce conformal transformations of the type of C_1 , given by (30) in order to have at the end the correct cycling transformations for the composite vertex. In order to do this we need to use matrices $E_{nm}(\gamma)$, defined by [5]

$$E_{nm}(\gamma) = \frac{1}{(m+1)!} \frac{\partial^{m+1}}{\partial z^{m+1}} \left[(\gamma z)^{n+1} \left(\frac{\partial}{\partial z} \gamma z \right)^{-1} \right] \Big|_{z=0} \quad (114)$$

which have the following properties:

$$\sum_{t=-1}^1 E_{rt}(\gamma_1) E_{ts}(\gamma_2) = E_{rs}(\gamma_1 \gamma_2) \quad , \quad r, s, t = -1, 0, 1 \quad , \quad (115)$$

$$E_{rn}(\gamma) = 0 \quad , \quad r = -1, 0, 1 \quad , \quad n \geq 2 \quad , \quad (116)$$

$$\sum_{p=-1}^{\infty} E_{np}(\gamma_1) E_{pm}(\gamma_2) = E_{nm}(\gamma_1 \gamma_2) \quad , \quad n, m \geq -1 \quad , \quad (117)$$

$$\sum_{p=2}^{\infty} E_{np}(\gamma_1) E_{pm}(\gamma_2) = E_{nm}(\gamma_1 \gamma_2) - \sum_{r,s=-1}^1 E_{nr}(\gamma_1) E_{rs}(\gamma_2) \delta_{sm} \quad , \quad n, m \geq 2 \quad . \quad (118)$$

The action of the operator C_1 on the ghosts b_n^i can then be calculated in the following way: first we write b_n^i in terms of an integral over the conformal operator $b_n^i(\xi_i)$

$$b_n^i = \oint_{\xi_{0i}=0} d\xi_{0i} (\xi_{0i})^{n+1} b^i(\xi_{0i}) \quad . \quad (119)$$

Then we insert the operator C_1 :

$$C_1 b_n^i C_1^{-1} = \oint_{\xi_{0i}=0} d\xi_{0i} (\xi_{0i})^{n+1} \left(\frac{d}{d\xi_{0i}} V_i^{-1} V_{0i} \xi_{0i} \right)^2 b^i(V_i^{-1} V_{0i} \xi_{0i}) \quad . \quad (120)$$

After a change of variables $\xi_i = V_i^{-1} V_{0i} \xi_{0i}$ we have⁷

$$C_1 b_n^i C_1^{-1} = \oint_{\xi_i=0} d\xi_i (V_{0i}^{-1} V_i \xi_i)^{n+1} \left(\frac{d}{d\xi_i} V_{0i}^{-1} V_i \xi_i \right)^{-1} b^i(\xi_i) \quad . \quad (121)$$

Using matrices $E_{nm}(\gamma)$, we then may expand ξ_{0i} in terms of ξ_i . If we also expand $b^i(\xi_i)$, we then obtain

$$C_1 b_n^i C_1^{-1} = \sum_{m=-1}^{\infty} \sum_{p=-\infty}^{\infty} \oint_{\xi_i=0} d\xi_i E_{nm}(V_{0i}^{-1} V_i) (\xi_i)^{m+1} b_{-p}^i(\xi_i)^{p-2} \quad . \quad (122)$$

Performing the integration we then have

$$C_1 b_n^i C_1^{-1} = \sum_{m=-1}^{\infty} E_{nm}(V_{0i}^{-1} V_i) b_m^i \quad . \quad (123)$$

Using (123) in equations (111) and (113) and multiplying (111) by $(d\xi_i/d\xi_{0i})^{-2}$ and (113) by $d\xi_i/d\xi_{0i}$, we then have (figure 14)

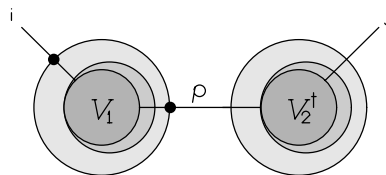


Figure 14: Overlap identity for $U_1^0 G C_1^{-1}$.

⁷Note that, because ξ_i is a polynomial in ξ_{0i} (with no constant term), then $\xi_{0i} = 0 \Rightarrow \xi_i = 0$.

$$U_1^0 G C_1^{-1} \left[b^i(\xi_i) - \left(\frac{d\xi_E}{d\xi_i} \right)^2 b^E(\xi_E) \right] = 0, \quad (124)$$

$$U_1^0 G C_1^{-1} \left[c^i(\xi_i) - \left(\frac{d\xi_E}{d\xi_i} \right)^{-1} c^E(\xi_E) \right] + V_1^0 C_1^{-1} \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \left[\sum_{l=1}^{N_1} \sum_{q, t=-1}^{\infty} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l \right] \\ \times \sum_{n, m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_E}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_E) (\xi_E)^{m+1} \right] = 0. \quad (125)$$

Before going further, some words must be said about the effects of C_1^{-1} on vertex U_1^0 with the extra ghosts G . This is given explicitly by [4][5]

$$U_1^0 G = \left(\prod_{i=1}^{N_1} i \langle 0 | \right) \exp \left[\sum_{\substack{i, j=1 \\ i \neq j}}^{N_1} \sum_{n=2}^{\infty} \sum_{m=-1}^{\infty} c_n^i E_{nm} (\Gamma V_{0i}^{-1} V_{0j}) b_m^j \right] \\ \times \prod_{r=-1}^1 \sum_{i=1}^{N_1} \sum_{s=-1}^1 E_{rs}(V_{0i}) b_s^i \times \prod_{\substack{i=1 \\ i \neq b, c}}^{N_1} \sum_{j=1}^N \sum_{n=-1}^{\infty} e_n^{ij} b_n^j. \quad (126)$$

Making use of matrices $F_{nm}(\gamma)$, defined by [5]

$$F_{nm}(\gamma) = \frac{1}{(m-2)!} \frac{\partial^{m-2}}{\partial z^{m-2}} \left\{ [\gamma(z)]^{n-2} \left[\frac{\partial}{\partial z} \gamma(z) \right]^{-1} \right\} \Big|_{z=0} \quad (127)$$

we may calculate in a similar way as we did for the b_n^i ghosts the effect of C_1 on the c_n^i ghosts, obtaining

$$C_1 c_n^i C_1^{-1} = \sum_{m=2}^{\infty} F_{nm}(V_{0i}^{-1} V_i) c_m^i. \quad (128)$$

Using this together with the property

$$F_{nm}(\gamma) = E_{mn}(\Gamma \gamma^{-1} \Gamma) \quad (129)$$

and equation (123), we may then show that the result of acting with C_1^{-1} on $U_1^0 G$ is

$$U_1^0 G C_1^{-1} = \left(\prod_{i=1}^{N_1} i \langle 0 | \right) \exp \left[\sum_{\substack{i, j=1 \\ i \neq j}}^{N_1} \sum_{n=2}^{\infty} \sum_{m=-1}^{\infty} c_n^i E_{nm} (\Gamma V_i^{-1} V_j) b_m^j \right] \\ \times \prod_{r=-1}^1 \sum_{i=1}^{N_1} \sum_{s=-1}^1 E_{rs}(V_i) b_s^i \times \prod_{\substack{i=1 \\ i \neq b, c}}^{N_1} \sum_{j=1}^N \sum_{n, m=-1}^{\infty} e_n^{ij} E_{nm}(V_{0j} V_j) b_m^j. \quad (130)$$

So we can see that in this case the action of C_1^{-1} on $U_1^0 G$ is not just to change $V_{0i} \rightarrow V_i$. Because of the peculiar nature of e_n^{ij} , it transforms as

$$e_n^{ij} \rightarrow \sum_{m=-1}^{\infty} e_n^{ij} E_{nm}(V_{0j} V_j) b_m^j. \quad (131)$$

Only in one particular group of cycling transformations (as we shall see later) will this be just equivalent to changing $V_{0i} \rightarrow V_i$. We shall call from now on $U_1^0 G C_1^{-1} \equiv U_1$ and $V_1^0 C_1^{-1} \equiv V_1$. The calculation for vertex $C_F U_2^\dagger C_2^{-1}$ will be similar to the one we have just made for $U_1^0 G C_1^{-1}$.

Having done this, we must insert the propagator \mathcal{P} into the overlap identities (124) and (125) in the same way as in the case with no ghosts. But now we must take extra care since there are terms depending on b_q^E in the second term of equation (125). Using equation (123) as a guideline, we have

$$\mathcal{P}^{-1}b_q^E\mathcal{P} = \sum_{t=-1}^{\infty} E_{qt}(\mathcal{P})b_t^E \quad (132)$$

so that the result of inserting \mathcal{P} into overlaps (124) and (125) is (figure 15)

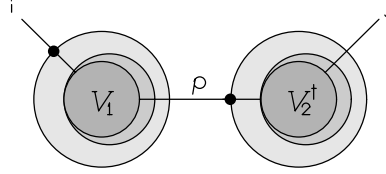


Figure 15: Overlap identity for $U_1\mathcal{P}$.

$$\begin{aligned} U_1\mathcal{P} \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^2 b^E(\mathcal{P}^{-1}\xi_E) \right] &= 0, \quad (133) \\ U_1\mathcal{P} \left[c^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^{-1} c^E(\mathcal{P}^{-1}\xi_E) \right] \\ &+ V_1\mathcal{P} \sum_{\substack{p=1 \\ p \neq b,c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_i)b_t^l + e_q^{kE} E_{qt}(V_{0E}^{-1}V_E\mathcal{P})b_t^E \right] \\ &\times \sum_{n,m=-1}^{\infty} \left[e_n^{pi} E_{nm}(V_{0i}^{-1}V_i)(\xi_i)^{m+1} - \left(\frac{d\xi_E}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm}(V_{0E}^{-1}V_E)(\xi_E)^{m+1} \right] = 0. \quad (134) \end{aligned}$$

The operators of the overlap equations are now facing leg F of vertex U_2^\dagger . In order to obtain the overlap identities for this leg, we must now identify the operators of leg E with the ones of leg F , which are adjoint operators:

$$b_t^E \rightarrow b_t^{F\dagger}, \quad b^E \rightarrow b^{F\dagger}, \quad c^E \rightarrow c^{F\dagger}. \quad (135)$$

First, as b^F and c^F are conformal operators with weights 2 and -1 , respectively, we have

$$b^{F\dagger}(\xi_F) = \Gamma b^F(\xi_F)\Gamma = \left(\frac{d}{d\xi_F} \Gamma\xi_F \right)^2 b^F(\Gamma\xi_F) = (\xi_F)^{-4} b^F(\Gamma\xi_F), \quad (136)$$

$$c^{F\dagger}(\xi_F) = \Gamma c^F(\xi_F)\Gamma = \left(\frac{d}{d\xi_F} \Gamma\xi_F \right)^{-1} c^F(\Gamma\xi_F) - (\xi_F)^2 c^F(\Gamma\xi_F). \quad (137)$$

Then, for b_n^F and c_n^F , we obtain

$$c_n^{F\dagger} = \Gamma c_n^F \Gamma = -c_{-n}^F, \quad (138)$$

$$b_n^{F\dagger} = \Gamma b_n^F \Gamma b_{-n}^F. \quad (139)$$

Then, we must also make the change

$$\begin{aligned} \left(\frac{d\xi_E}{d\xi_i} \right)^{-1} e_n^{pE}(\xi_E)^{n+1} &= \left(\frac{d\xi_F}{d\xi_i} \right)^{-1} \left(\frac{d\xi_E}{d\xi_F} \right)^{-1} e_n^{pE}(\xi_E)^{n+1} \\ &= \left(\frac{d\xi_F}{d\xi_i} \right)^{-1} \sum_{m=-1}^{\infty} e_n^{pE} E_{nm}(V_E^{-1}V_F)(\xi_F)^{m+1}. \quad (140) \end{aligned}$$

So, the overlap equations become

$$U_1 \mathcal{P} \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^2 (\mathcal{P}^{-1} \xi_E)^{-4} b^F(\Gamma \mathcal{P}^{-1} \xi_E) \right] = 0, \quad (141)$$

$$U_1 \mathcal{P} \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^{-1} (\mathcal{P}^{-1} \xi_E)^2 c^F(\Gamma \mathcal{P}^{-1} \xi_E) \right] \\ + V_1 \mathcal{P} \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \sum_{q, t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P}) b_{-t}^F \right] \\ \times \sum_{n, m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_F}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_F) (\xi_F)^{m+1} \right] = 0. \quad (142)$$

We are then facing the conformal transformation C_F that takes ξ_F into ξ_{0F} . Inserting this transformation we obtain (figure 16):

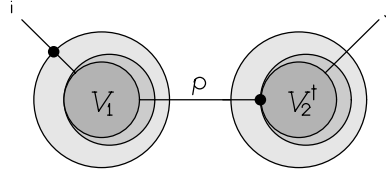


Figure 16: Overlap identity for $U_1 \mathcal{P} C_F$.

$$U_1 \mathcal{P} C_F \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^2 (\mathcal{P}^{-1} \xi_E)^{-4} \right. \\ \left. \times \left(\frac{dV_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^2 b^F(V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0, \quad (143)$$

$$U_1 \mathcal{P} C_F \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^{-1} (\mathcal{P}^{-1} \xi_E)^2 \left(\frac{dV_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^{-1} c^F(V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] \\ + V_1 \mathcal{P} C_F \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \sum_{q, t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l - e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_{0F}) b_t^F \right] \\ \times \sum_{n, m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_{0F}}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_{0F}) (\xi_{0F})^{m+1} \right] = 0. \quad (144)$$

We are facing now vertex $V_2^{0\dagger}$. This vertex satisfies the following overlap identity [4]:

$$\sum_{s=-1}^1 E_{rs}(V_{0F}) b_{-s}^F V_2^{0\dagger} = -V_2^{0\dagger} \sum_{\substack{i=1 \\ i \neq F}}^{N_2} \sum_{s=-1}^1 E_{rs}(V_{0i}) b_s^i, \quad r = -1, 0, 1. \quad (145)$$

Using equation (117), we have

$$E_{qt}(V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_{0F}) b_t^F = \sum_{u=-1}^{\infty} E_{qu}(V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1}) E_{ut}(V_{0F}) b_t^F. \quad (146)$$

All terms on b_t^F with $t \geq 2$ get annihilated by the conjugate vacuum $|0\rangle_F$, while we may use identity (145) to substitute the terms in b_r^F , $r = -1, 0, 1$. Doing this, equation (144) becomes

$$\begin{aligned}
& U_1 \mathcal{P} C_F \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^{-1} (\mathcal{P}^{-1} \xi_E)^2 \left(\frac{dV_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^{-1} c^F (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] \\
& + V_1 \mathcal{P} C_F \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \sum_{q, t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_{0l}) b_t^l \right] \\
& \times \sum_{n, m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_{0F}}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_{0F}) (\xi_{0F})^{m+1} \right] = 0. \quad (147)
\end{aligned}$$

Equations (143) and (147) are the overlap identities between legs i and F .

In order to obtain the overlaps between leg i of vertex U_1 and an arbitrary leg j of vertex U_2^\dagger , we must now perform a cycling transformation that will take the operators from leg F to leg j . The effect of this transformation on $c^F (\Gamma \mathcal{P}^{-1} V_E^{-1} V_i \xi_i)$ is

$$V_j^{-1} V_F c^F (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) V_F^{-1} V_j = \left(\frac{dV_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{dV_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E} \right)^{-1} c^j (V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E). \quad (148)$$

Then, we must also write

$$\left(\frac{d\xi_{0F}}{d\xi_i} \right)^{-1} \sum_{m=-1}^{\infty} e_n^{pE} E_{nm} (V_{0E}^{-1} V_{0F}) (\xi_{0F})^{m+1} = \left(\frac{d\xi_{0j}}{d\xi_i} \right)^{-1} \sum_{m=-1}^{\infty} e_n^{pE} E_{nm} (V_{0E}^{-1} V_{0j}) (\xi_{0j})^{m+1}. \quad (149)$$

Doing this, overlap equations (143, 147) become (figure 17)

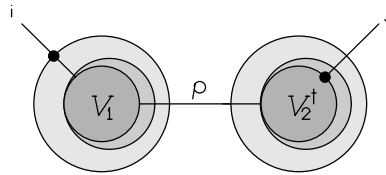


Figure 17: Overlap identity for $U_1 \mathcal{P} C_F V_2^{0\dagger}$.

$$\begin{aligned}
& U_1 \mathcal{P} C_F V_2^{0\dagger} \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^2 (\mathcal{P}^{-1} \xi_E)^{-4} \right. \\
& \quad \left. \times \left(\frac{dV_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^2 b^F (V_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] = 0, \quad (150)
\end{aligned}$$

$$\begin{aligned}
& U_1 \mathcal{P} C_F V_2^{0\dagger} \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} \xi_E \right)^{-1} (\mathcal{P}^{-1} \xi_E)^2 \left(\frac{dV_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^{-1} c^F (V_{0j}^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E) \right] \\
& + V_1 \mathcal{P} C_F V_2^{0\dagger} \sum_{\substack{p=1 \\ p \neq b, c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b, c \\ k \neq p}}^{N_1} \sum_{q, t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_{0l}) b_t^l \right] \\
& \times \sum_{n, m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_{0j}}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_{0j}) (\xi_{0j})^{m+1} \right] = 0. \quad (151)
\end{aligned}$$

The operators are now facing the ghosts that surround vertex $U_2^{0\dagger}$ (like in equation (74)):

$$\prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left(\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{q=-1}^{\infty} e_q^{kF} b_{-q}^F \right). \quad (152)$$

so that we must insert these ghosts into the expressions for the overlaps. Before doing that, we must notice that the extra ghosts acting on vertex $V_2^{0\dagger}$ have at their left both the conformal transformation C_F and the propagator \mathcal{P} so that we must first pass them through in order to reach vertex U_1 :

$$\begin{aligned} & \mathcal{P} C_F \prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left(\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{q=-1}^{\infty} e_q^{kF} b_{-q}^F \right) \\ &= \prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left(\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l - \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1}) b_t^F \right) \mathcal{P} C_F. \end{aligned} \quad (153)$$

Then, identifying legs E and F , we have the following expression for the ghosts:

$$\prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l - \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1}) b_t^E \right] \mathcal{P}. \quad (154)$$

We may now pass it through the conformal transformation C_1 , obtaining

$$\prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l - \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_{0E}) b_t^E \right] \mathcal{P}. \quad (155)$$

This is now facing vertex V_1^0 , which satisfies the following overlap identity [4]:

$$V_1^0 \sum_{s=-1}^1 E_{rs} (V_{0E}) b_s^E = -V_1^0 \sum_{\substack{i=1 \\ i \neq E}}^{N_1} \sum_{s=-1}^1 E_{rs} (V_{0i}) b_s^i, \quad r = -1, 0, 1. \quad (156)$$

Using this identity, we may then write the extra ghosts as

$$\prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_{0l}) b_t^l \right]. \quad (157)$$

Passing it back through C_1 we then have

$$\prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_l) b_t^l \right]. \quad (158)$$

We now insert these ghosts into expressions (150) and (151). We do so by multiplying them by (158) and passing it through the first term of the overlaps. What we obtain is (figure 18)

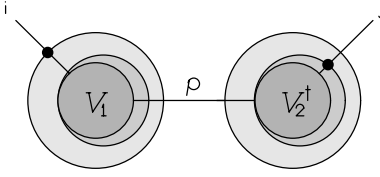


Figure 18: Overlap identity for $U_1\mathcal{P}C_FU_2^{0\dagger}$.

$$\begin{aligned}
& U_1\mathcal{P}C_FU_2^{0\dagger} \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^2 (\mathcal{P}^{-1}\xi_E)^{-4} \right. \\
& \quad \left. \times \left(\frac{dV_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^2 b^F(V_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E) \right] = 0, \quad (159) \\
& U_1\mathcal{P}C_FU_2^{0\dagger} \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^{-1} (\mathcal{P}^{-1}\xi_E)^2 \left(\frac{dV_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^{-1} c^F(V_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E) \right] \\
& + U_1\mathcal{P}C_FV_2^{0\dagger} \sum_{p=1}^{N_2} (-1)^p \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_l) b_t^l \right] \\
& \times \sum_{n=-1}^{\infty} \left[\sum_{m=-1}^{\infty} e_n^{pF} E_{nm} (V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_i) (\xi_i)^{m+1} \right. \\
& \quad \left. + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^{-1} (\mathcal{P}^{-1}\xi_E)^2 \left(\frac{dV_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^2 e_n^{pj} (V_{0j}^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E)^{n+1} \right] \\
& + V_1\mathcal{P}C_FV_2^{0\dagger} \sum_{\substack{p=1 \\ p \neq b,c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1}V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1}V_E\mathcal{P}\Gamma V_F^{-1}V_{0l}) b_t^l \right] \\
& \times \sum_{n,m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1}V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_{0j}}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1}V_E) (\xi_{0j})^{m+1} \right] \\
& \times \prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} \sum_{q=-1}^{\infty} e_q^{kl} b_q^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} \sum_{q,t=-1}^{\infty} e_q^{kF} E_{qt} (V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_l) b_t^l \right] = 0. \quad (160)
\end{aligned}$$

We are then facing the last term of this composite vertex: the conformal transformation C_2 on vertex $V_2^{0\dagger}$. Inserting it into equations (159) and (160), we obtain (figure 19)

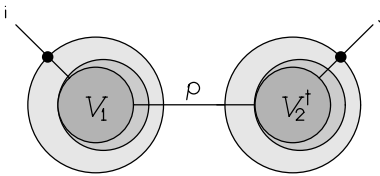


Figure 19: Overlap identity for $U_1\mathcal{P}U_2^\dagger$.

$$U_1\mathcal{P}U_2^\dagger \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i} \mathcal{P}^{-1}\xi_E \right)^2 (\mathcal{P}^{-1}\xi_E)^{-4} \right.$$

$$\times \left(\frac{dV_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^2 b^F \left(V_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E \right) \Big] = 0, \quad (161)$$

$$\begin{aligned} & U_1\mathcal{P}U_2^\dagger \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i}\mathcal{P}^{-1}\xi_E \right)^{-1} (\mathcal{P}^{-1}\xi_E)^2 \left(\frac{dV_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^{-1} c^F \left(V_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E \right) \right] \\ & + U_1\mathcal{P}V_2^\dagger \sum_{p=1}^{N_2} (-1)^p \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} \left(V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_l \right) b_t^l \right] \\ & \times \sum_{n=-1}^{\infty} \left[\sum_{m=-1}^{\infty} e_n^{pF} E_{nm} \left(V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_i \right) (\xi_i)^{m+1} \right. \\ & \quad \left. + \left(\frac{d}{d\xi_i}\mathcal{P}^{-1}\xi_E \right)^{-1} (\mathcal{P}^{-1}\xi_E)^2 \left(\frac{dV_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E}{d\Gamma\mathcal{P}^{-1}\xi_E} \right)^2 e_n^{pj} \left(V_j^{-1}V_F\Gamma\mathcal{P}^{-1}\xi_E \right)^{n+1} \right] \\ & + V_1\mathcal{P}V_2^\dagger \sum_{\substack{p=1 \\ p \neq b,c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} \left(V_{0E}^{-1}V_E\mathcal{P}\Gamma V_F^{-1}V_l \right) b_t^l \right] \\ & \times \sum_{n,m=-1}^{\infty} \left[e_n^{pi} E_{nm}(V_{0i}^{-1}V_i)(\xi_i)^{m+1} - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm}(V_{0E}^{-1}V_E)(\xi_j)^{m+1} \right] \\ & \times \prod_{\substack{k=1 \\ k \neq d,g,h}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + \sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} \left(V_{0F}^{-1}V_F\Gamma\mathcal{P}^{-1}V_E^{-1}V_l \right) b_t^l \right] = 0 \end{aligned} \quad (162)$$

where we have called $C_F U_2^{0\dagger} C_2^{-1} \equiv U_2^\dagger$ and $C_F V_2^{0\dagger} C_2^{-1} \equiv V_2^\dagger$.

We must now extract the ghosts from vertex U_1 in the second term of equation (162):

$$\begin{aligned} U_1\mathcal{P} &= V_1 \prod_{\substack{k=1 \\ k \neq b,c}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + e_q^{kE} E_{qt}(V_{0E}^{-1}V_E)b_t^E \right] \mathcal{P} \\ &= V_1\mathcal{P} \prod_{\substack{k=1 \\ k \neq b,c}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + e_q^{kE} E_{qt}(V_{0E}^{-1}V_E\mathcal{P})b_t^E \right]. \end{aligned} \quad (163)$$

Identifying b_t^E with b_{-t}^F and passing now these ghosts through the conformal transformation C_F , we obtain

$$U_1\mathcal{P}C_F = V_1\mathcal{P}C_F \prod_{\substack{k=1 \\ k \neq b,c}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l - e_q^{kE} E_{qt}(V_{0E}^{-1}V_E\mathcal{P}\Gamma V_F^{-1}V_{0F})b_t^F \right]. \quad (164)$$

Using now overlap identities (145) for vertex V_2^\dagger and inserting the conformal transformation C_2 , we then obtain

$$U_1\mathcal{P}C_F V_2^\dagger = V_1\mathcal{P}C_F V_2^\dagger \prod_{\substack{k=1 \\ k \neq b,c}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq F}}^{N_1} e_q^{kl} E_{qt}(V_{0l}^{-1}V_l)b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} \left(V_{0E}^{-1}V_E\mathcal{P}\Gamma V_F^{-1}V_l \right) b_t^l \right]. \quad (165)$$

Substituting (165) into equation (162) we then obtain the overlap identities between legs i and j :

$$U_c \left[b^i(\xi_i) - \left(\frac{d}{d\xi_i}\mathcal{P}^{-1}V_E^{-1}V_i\xi_i \right)^2 (\mathcal{P}^{-1}V_E^{-1}V_i\xi_i)^{-4} \right]$$

$$\begin{aligned}
& \times \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^2 b^j \left(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right) \Big] = 0, \quad (166) \\
U_c & \left[c^i(\xi_i) + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^{-1} \left(\mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^2 \right. \\
& \left. \times \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^{-1} c^j \left(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right) \right] \\
& + V_c \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_l) b_t^l \right] \\
& \times \sum_{\substack{p=1 \\ p \neq d,g,h}}^{N_2} (-1)^p \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l \right] \\
& \times \sum_{n,m=-1}^{\infty} \left[e_q^{pF} E_{nm} (V_{0F}^{-1} V_F \mathcal{P}^{-1} V_E^{-1} V_i) (\xi_i)^{m+1} + \left(\frac{d}{d\xi_i} \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^{-1} \left(\mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^2 \right. \\
& \left. \times \left(\frac{dV_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E}{d\Gamma \mathcal{P}^{-1} \xi_E} \right)^2 e_n^{pj} E_{nm} (V_{0j}^{-1} V_j) \left(V_j^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_i \xi_i \right)^{m+1} \right] \\
& + V_c \sum_{\substack{p=1 \\ p \neq b,c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_E \mathcal{P} \Gamma V_F^{-1} V_l) b_t^l \right] \\
& \times \sum_{n,m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_j) (\xi_j)^{n+1} \right] \\
& \times \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} (V_{0F}^{-1} V_F \Gamma \mathcal{P}^{-1} V_E^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l \right] = 0. \quad (167)
\end{aligned}$$

If we now impose that these are the correct overlap equations between legs i and j of the composite vertex we then must have:

$$V_j^{-1} V_F \Gamma \mathcal{P}^{-1} \xi_E = \xi_j \quad (168)$$

what fixes the propagator as

$$\mathcal{P} = V_E^{-1} V_F \Gamma \quad (169)$$

which is the same form of the propagator for the bosonic part, but now with the cyclings defined on the complete generators \mathcal{L}_n^i . The overlap equations now read

$$U_c \left[b^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^2 b^j(\xi_j) \right] = 0, \quad (170)$$

$$\begin{aligned}
U_c & \left[c^i(\xi_i) - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} c^j(\xi_j) \right] + V_c \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_l) b_t^l \right] \\
& \times \sum_{\substack{p=1 \\ p \neq d,g,h}}^{N_2} (-1)^p \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} (V_{0F}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l \right]
\end{aligned}$$

$$\begin{aligned}
& \times \sum_{n,m=-1}^{\infty} \left[e_n^{pF} E_{nm} (V_{0F}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} e_n^{pj} E_{nm} (V_{0j}^{-1} V_j) (\xi_j)^{m+1} \right] \\
& + V_c \sum_{\substack{p=1 \\ p \neq b,c}}^{N_1} (-1)^p \prod_{\substack{k=1 \\ k \neq b,c \\ k \neq p}}^{N_1} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kE} E_{qt} (V_{0E}^{-1} V_l) b_t^l \right] \\
& \times \sum_{n,m=-1}^{\infty} \left[e_n^{pi} E_{nm} (V_{0i}^{-1} V_i) (\xi_i)^{m+1} - \left(\frac{d\xi_j}{d\xi_i} \right)^{-1} e_n^{pE} E_{nm} (V_{0E}^{-1} V_j) (\xi_j)^{n+1} \right] \\
& \times \prod_{\substack{k=1 \\ k \neq d,g,h \\ k \neq p}}^{N_2} \sum_{q,t=-1}^{\infty} \left[\sum_{\substack{l=1 \\ l \neq E}}^{N_1} e_q^{kF} E_{qt} (V_{0F}^{-1} V_l) b_t^l + \sum_{\substack{l=1 \\ l \neq F}}^{N_2} e_q^{kl} E_{qt} (V_{0l}^{-1} V_l) b_t^l \right] = 0 . \tag{171}
\end{aligned}$$

These are the overlap identities between legs i and j of the composite vertex U_c .

From these overlap equations it is then possible to derive the form of the composite vertex. It is given by

$$\begin{aligned}
U_c^{\text{gh}} &= \left(\prod_{\substack{i=1 \\ i \neq E,F}}^{N_1+N_2} i \langle 0| \right) \exp \left[\sum_{\substack{i,j=1 \\ i \neq j \\ i,j \neq E,F}}^{N_1+N_2} \sum_{n=2}^{\infty} \sum_{m=-1}^{\infty} c_n^{ij} E_{nm} (\Gamma V_i^{-1} V_j) b_m^j \right] \times \prod_{r=-1}^1 \sum_{\substack{i=1 \\ i \neq E,F}}^{N_1+N_2} \sum_{s=-1}^1 E_{rs} (V_i) b_s^i \\
& \times \prod_{\substack{i=1 \\ i \neq b,c}}^{N_1} \sum_{n,m=-1}^{\infty} \left[\sum_{\substack{j=1 \\ j \neq E}}^{N_1} e_n^{ij} E_{nm} (V_{0j}^{-1} V_j) b_m^j + \sum_{\substack{j=1 \\ j \neq F}}^{N_2} e_n^{iE} E_{mn} (V_{0E}^{-1} V_j) b_n^j \right] \\
& \times \prod_{\substack{i=1 \\ i \neq d,g,h}}^{N_2} \sum_{n,m=-1}^{\infty} \left[\sum_{\substack{j=1 \\ j \neq E}}^{N_1} e_n^{iF} E_{nm} (V_{0F}^{-1} V_j) b_n^j + \sum_{\substack{j=1 \\ j \neq F}}^{N_2} e_n^{ij} E_{nm} (V_{0j}^{-1} V_j) b_m^j \right] . \tag{172}
\end{aligned}$$

Although this is the correct composite vertex for a general cycling, its ghost structure is not very apparent. We may use the explicit expression for the vectors e_n^{ij} [7]:

$$e_n^{ij} = \sum_{m=-1}^{\infty} k_m^{ij} E_{mn} (\gamma_m^j) \tag{173}$$

where

$$k_{-1}^{ij} = \delta_{ij} \quad , \quad k_0^{ij} = -\frac{\partial}{\partial z_i} \ln a_0^j \quad , \quad k_n^{ij} = -\frac{\partial}{\partial z_i} \quad , \quad n \geq 1 \tag{174}$$

and

$$\gamma_{-1}^j = V_j \quad , \tag{175}$$

$$\gamma_0^j = \exp(-\mathcal{L}_0^j \ln a_0^j) \exp\left(-\sum_{n=1}^{\infty} \bar{a}_n^j \mathcal{L}_n^j\right) \quad , \tag{176}$$

$$\gamma_p^j = \exp\left(-\sum_{n=p+1}^{\infty} \bar{a}_n^j \mathcal{L}_n^j\right) \quad , \quad p \geq 1 . \tag{177}$$

If we assume now that the cycling transformations V_E^{-1} and V_F^{-1} do not depend on the variables z_i of vertex U_1 or z_j of vertex U_2^\dagger , and if we have in mind that $e_n^{ij} = 0$ for any leg i of vertex U_1 and a leg j of vertex U_2^\dagger (or vice-versa), then we have

$$e_n^{ij} = \delta^{ij} \sum_{r=-1}^1 E_{-1r} (V_i) \tag{178}$$

what simplifies things considerably. We shall call all cyclings that have such properties “simple cycling-like”. In this kind of cyclings, each leg has its own ghost attached to it, with the exception of three of the legs which have no ghosts attached to them. We will now consider three cases separately: one in which none of the legs E or F have ghosts attached to them, one in which one of these legs (say E) has a ghost attached to it, and a case where both legs (E and F) have ghosts attached to them.

The composite vertex for the case where neither E nor F (we choose $E = b$ and $F = g$) have ghosts attached is given by

$$U_c^{\text{gh}} = V_c^{\text{gh}} \times \prod_{\substack{i=1 \\ i \neq b \\ i \neq E}}^{N_1} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i \times \prod_{\substack{i=1 \\ i \neq d, h \\ i \neq F}}^{N_2} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i \quad (179)$$

where V_c^{gh} is the vertex given by the first two terms of equation (172). For the case where leg E has a ghost attached to it, but not leg F (we shall call $F = g$), the composite vertex is given by

$$U_c^{\text{gh}} = V_c^{\text{gh}} \times \prod_{\substack{i=1 \\ i \neq b, c \\ i \neq E}}^{N_1} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i \times \sum_{\substack{j=1 \\ j \neq F}}^{N_2} \sum_{r=-1}^1 E_{-1r}(V_j) b_r^j \times \prod_{\substack{i=1 \\ i \neq d, h \\ i \neq F}}^{N_2} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i . \quad (180)$$

In the last case, where both E and F have ghosts attached to them, we then have

$$\begin{aligned} U_c^{\text{gh}} &= V_c^{\text{gh}} \times \prod_{\substack{i=1 \\ i \neq b, c \\ i \neq E}}^{N_1} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i \\ &\times \sum_{\substack{j=1 \\ j \neq F}}^{N_2} \sum_{r=-1}^1 E_{-1r}(V_j) b_r^j \times \sum_{\substack{j=1 \\ j \neq E}}^{N_1} \sum_{r=-1}^1 E_{-1r}(V_j) b_r^j \times \prod_{\substack{i=1 \\ i \neq d, g, h \\ i \neq F}}^{N_2} \sum_{r=-1}^1 E_{-1r}(V_i) b_r^i . \end{aligned} \quad (181)$$

We can see that, for each case, the composite vertex has the correct ghosts number ($N_1 + N_2 - 2$).

4 Conclusions

Using overlap identities, two vertices were sewn together in order to become a composite vertex. The calculations have been done with the correct ghost numbers for each vertex and the result has both BRST invariance and the correct ghost counting.

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References

- [1] A. Neveu and P. West, Cycling, twisting and sewing in the Group Theoretic Approach to Strings, *Comm.Math.Phys.* 119 (1988) 585.
- [2] A. Neveu and P. West, Group Theoretic Approach to the open bosonic string multi-loop S-matrix, *Comm.Math.Phys.* 114 (1988) 613.
- [3] P. Di Vecchia, F. Pezzella, M. Frau, K. Hornfeck, A. Lerda and S. Sciuto, N -point g -loop vertex for the free bosonic theory with vacuum charge Q , *Nucl.Phys.* B322 (1989) 317.

- [4] M. Freeman and P. West, Ghost vertices for the bosonic string using the Group Theoretic approach to String Theory, Phys.Lett. 205B (1988) 30.
- [5] P. Di Vecchia, F. Pezzella, M. Frau, K. Hornfeck, A. Lerda and S. Sciuto, N -point g -loop vertex for a free fermionic theory with arbitrary spin, Nucl.Phys. B333 (1990) 635
- [6] P. Di Vecchia, M. Frau, A. Lerda and S. Sciuto, A simple expression for the multiloop amplitude in the bosonic string, Phys. Lett. 199B (1987) 49
P. Di Vecchia, M. Frau, A. Lerda and S. Sciuto, n -string vertex and loop calculation in the bosonic string, Nucl. Phys. B298 (1988) 527
- [7] L. Sandoval Jr., String Scattering, Ph.D. thesis (1995) King's College London, Univ. of London (excerpts available with the author).